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## Thermal stability of $\text{Ba}(\text{Zr}_{0.8-x}\text{Ce}_x\text{Y}_{0.2})\text{O}_{2.9}$ ceramics in carbon dioxide

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*In situ* x-ray diffraction spectra (25–1000 °C) have been measured as a function of temperature for proton-conducting  $\text{Ba}(\text{Zr}_{0.8-x}\text{Ce}_x\text{Y}_{0.2})\text{O}_{2.9}$  ( $x=0.0-0.4$ ) ceramics in  $\text{CO}_2$  atmosphere. Atomic vibrations before and after exposure to  $\text{CO}_2$  were obtained by using the micro-Raman scattering (150–1600  $\text{cm}^{-1}$ ).  $\text{Ba}(\text{Zr}_{0.8}\text{Y}_{0.2})\text{O}_{2.9}$  and  $\text{Ba}(\text{Zr}_{0.6}\text{Ce}_{0.2}\text{Y}_{0.2})\text{O}_{2.9}$  reveal a promising thermal stability in  $\text{CO}_2$  without apparent decomposition up to 1000 °C. However,  $\text{Ba}(\text{Zr}_{0.5}\text{Ce}_{0.3}\text{Y}_{0.2})\text{O}_{2.9}$  and  $\text{Ba}(\text{Zr}_{0.4}\text{Ce}_{0.4}\text{Y}_{0.2})\text{O}_{2.9}$  exhibit thermally stable below 550 °C and then proceed an obvious chemical decomposition of  $\text{BaCO}_3$  and  $\text{Zr}_{0.8-x}\text{Ce}_x\text{Y}_{0.2}\text{O}_2$  above 550 °C, which were clearly evidenced by the Raman vibrations of 1057 and 466  $\text{cm}^{-1}$ , respectively. A first-order orthorhombic-hexagonal structure transition was confirmed in  $\text{BaCO}_3$  in the region of 810–850 °C upon heating. This study suggests that the  $\text{Ba}(\text{Zr}_{0.8-x}\text{Ce}_x\text{Y}_{0.2})\text{O}_{2.9}$  ceramics with  $x \leq 0.2$  are promising candidates for proton-conducting applications in  $\text{CO}_2$ -containing environment. © 2009 American Institute of Physics. [DOI: 10.1063/1.3117835]

### I. INTRODUCTION

One challenging issue in proton-conducting ceramics for applications of hydrogen purification and solid oxide fuel cell (SOFC) electrolytes is thermal instability due to the reaction with environment, especially components of hydrocarbon fuels such as  $\text{CO}_2$ ,  $\text{H}_2\text{S}$ , and other trace species.<sup>1-3</sup> Such unfavorable reactions with carbon species and  $\text{H}_2\text{S}$  could cause performance loss and degradation. However, the effect of species of hydrocarbon fuels on SOFC materials is not presently well known. Although proton conductors are promising candidates for SOFC at intermediate temperatures (typically 700–850 °C) because of their low activation energy,<sup>4-7</sup> the remained challenge is to find an appropriate compromise between ionic conductivity and thermal stability in various environments.

Doped  $\text{BaCeO}_3$  has been known to exhibit high ionic conductivity above 500 °C and poor stability in  $\text{CO}_2$  and  $\text{H}_2\text{O}$ .<sup>4-13</sup> On the other hand, yttrium-doped  $\text{BaZrO}_3$  shows a sufficient thermal stability.<sup>14,15</sup> Partially substituting Zr for Ce can reduce tendency of decomposition in  $\text{CO}_2$  at high temperature but decreases the ionic conductivity. Therefore, it has been a goal to find doped  $\text{Ba}(\text{Zr},\text{Ce})\text{O}_3$  ceramics with sufficient ionic conductivity and thermal stability by replacing a fraction of cerium with Zr or other dopants.<sup>14-22</sup>

By thermal gravimetric analysis (TGA) and x-ray diffraction (XRD),  $\text{BaCe}_{0.9}\text{Y}_{0.1}\text{O}_{3-\delta}$  (BCY10) powder was confirmed to be only kinetically stable below 500 °C and then decomposes completely to  $\text{BaCO}_3$ ,  $\text{CeO}_2$ , and  $\text{Y}_2\text{O}_3$  after heating in pure  $\text{CO}_2$  at 860 °C.<sup>8</sup> Similarly, powder XRD result of BCY10 heated in 100%  $\text{CO}_2$  in the range of 700–1000 °C evidenced  $\text{BaCO}_3$  and fluoritelike  $\text{CeO}_2$

structure.<sup>9</sup> It was found that BCY10 can absorb 0.13 g of  $\text{CO}_2$  per ceramic gram.<sup>9</sup>  $\text{BaCe}_{0.9}\text{Nd}_{0.1}\text{O}_{3-\delta}$  (BCN) ceramic also showed decomposition of  $\text{BaCO}_3$  and  $\text{CeO}_2$  in 1 atm  $\text{CO}_2$  before reaching 1200 °C.<sup>10</sup> Above 1200 °C, the BCN ceramics reacted with alumina or zirconia, leading to the loss of barium and an excess of cerium.<sup>10</sup>

The XRD result of  $\text{BaZr}_{0.4}\text{Ce}_{0.5}\text{Y}_{0.1}\text{O}_{3-\delta}$  showed a good stability after exposure to  $\text{CO}_2$  at 900 °C.<sup>12</sup> The total conductivity of  $\text{BaZr}_x\text{Ce}_{0.9-x}\text{Y}_{0.1}\text{O}_{3-\delta}$  ( $x=0.0-0.9$ ) in wet  $\text{H}_2$  ( $p\text{H}_2\text{O}=1.7 \times 10^3$  Pa) varies from  $3.5 \times 10^{-2}$  to  $3.0 \times 10^{-3}$   $\text{S cm}^{-1}$  at 800 °C.<sup>12</sup> The XRD spectra of  $\text{Ba}(\text{Zr}_{0.4}\text{Ce}_{0.5}\text{Y}_{0.1})\text{O}_{2.95}$  and  $\text{Ba}(\text{Zr}_{0.6}\text{Ce}_{0.3}\text{Y}_{0.1})\text{O}_{2.95}$  sintered pellets showed good stability after being boiled in water or after being exposed to  $\text{CO}_2$  at 900 °C.<sup>15</sup> TGA and differential thermal analysis of Gd- and Nd-doped  $\text{Ba}(\text{Ce},\text{Zr})\text{O}_3$  solid solutions showed a reaction with  $\text{CO}_2$  above 600 °C and a reverse reaction at about 1150 °C for low Zr content.<sup>16</sup>

Recent XRD investigation of  $\text{Ba}(\text{Ce}_{0.8}\text{Y}_{0.2})\text{O}_3$  and  $\text{Ba}(\text{Zr}_{0.1}\text{Ce}_{0.7}\text{Y}_{0.2})\text{O}_{3-\delta}$  powders exposed to 2%  $\text{CO}_2$  (with  $\text{H}_2$ ) at 500 °C for 1 week suggested that  $\text{Ba}(\text{Ce}_{0.8}\text{Y}_{0.2})\text{O}_3$  decomposed to  $\text{BaCO}_3$ ,  $\text{CeO}_2$ , and  $\text{Y}_2\text{O}_3$ .<sup>18</sup> However,  $\text{Ba}(\text{Zr}_{0.1}\text{Ce}_{0.7}\text{Y}_{0.2})\text{O}_{3-\delta}$  remained unchanged at 500 °C and exhibited sufficient stability in 2%  $\text{CO}_2$  atmosphere.<sup>18</sup> The XRD result of  $\text{Ba}(\text{Zr}_{0.1}\text{Ce}_{0.7}\text{Y}_{0.2})\text{O}_{3-\delta}$  powder before and after exposure to  $\text{H}_2$  containing 15%  $\text{H}_2\text{O}$  at 500 °C showed no decomposition.<sup>18</sup> Most recently, double perovskitelike  $\text{Ba}_2(\text{Ca}_{0.75}\text{Nb}_{0.59}\text{Ta}_{0.66})\text{O}_{6-\delta}$ ,  $\text{Ba}_2(\text{Ca}_{0.75}\text{Nb}_{0.66}\text{Ta}_{0.59})\text{O}_{6-\delta}$  and  $\text{Ba}_2(\text{Ca}_{0.79}\text{Nb}_{0.66}\text{Ta}_{0.55})\text{O}_{6-\delta}$  showed a long-term structural stability in  $\text{CO}_2$  at 800 °C and boiling water.<sup>23</sup>

Although proton-conducting ceramics are promising candidates for SOFC, the major issue for these materials is to find appropriate components with a wide-temperature-range thermal stability in various environments. In this work, *in situ* temperature-dependent XRD and post-micro-Raman scattering were employed to investigate the structure stability

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of  $\text{Ba}(\text{Zr}_{0.8-x}\text{Ce}_x\text{Y}_{0.2})\text{O}_{2.9}$  ( $x=0.0, 0.2, 0.3,$  and  $0.4$ ) ceramic powders before and after exposure to  $\text{CO}_2$  in the temperature region of  $25\text{--}1000\text{ }^\circ\text{C}$ .

## II. EXPERIMENTAL PROCEDURE

$\text{Ba}(\text{Zr}_{0.8-x}\text{Ce}_x\text{Y}_{0.2})\text{O}_{2.9}$  ( $x=0.0, 0.2, 0.3,$  and  $0.4$ ) ceramic powders were synthesized by the glycine-nitrate process.<sup>24</sup> Synthesized powders were calcined at  $1300\text{ }^\circ\text{C}$  for 5 h to optimize the perovskite phase. Hereafter, BZCY442, BZCY532, BZCY622, and BZY82 represent  $\text{Ba}(\text{Zr}_{0.4}\text{Ce}_{0.4}\text{Y}_{0.2})\text{O}_{2.9}$ ,  $\text{Ba}(\text{Zr}_{0.5}\text{Ce}_{0.3}\text{Y}_{0.2})\text{O}_{2.9}$ ,  $\text{Ba}(\text{Zr}_{0.6}\text{Ce}_{0.2}\text{Y}_{0.2})\text{O}_{2.9}$ , and  $\text{Ba}(\text{Zr}_{0.8}\text{Y}_{0.2})\text{O}_{2.9}$ , respectively.  $\text{Ba}(\text{Zr,Ce,Y})\text{O}_3$  and  $(\text{Zr,Ce,Y})\text{O}_2$  symbolize  $\text{Ba}(\text{Zr}_{0.8-x}\text{Ce}_x\text{Y}_{0.2})\text{O}_{2.9}$  and  $\text{Zr}_{0.8-x}\text{Ce}_x\text{Y}_{0.2}\text{O}_2$ , respectively. For *in situ* XRD measurements, a Rigaku Model MultiFlex x-ray diffractometer with  $\text{Cu K}\alpha_1$  and  $\text{Cu K}\alpha_2$  radiations was used. The intensity ratio between  $\text{K}\alpha_1$  and  $\text{K}\alpha_2$  is about 2:1.<sup>25</sup> The calcined powders were scanned at room temperature first before flowing  $\text{CO}_2$ , and then the atmosphere was switched to flow at 1 atm  $\text{CO}_2$ . The temperature was raised in steps from room temperature. Each XRD scan was taken after holding the powder for more than 40 min at the setting temperature to allow complete reaction with  $\text{CO}_2$ .

A double grating Jobin Yvon Model U-1000 double monochromator with 1800 grooves/mm gratings and a nitrogen-cooled charge coupled device as a detector were employed for the post-micro-Raman scattering. A Coherent Model Innova 90 argon laser with wavelength  $\lambda = 514.5\text{ nm}$  was used as an excitation source. The Raman scattering was performed in the region of  $150\text{--}1600\text{ cm}^{-1}$ . The powder used in the post-Raman scattering measurements was previously exposed to  $\text{CO}_2$  during the *in situ* XRD.

## III. RESULTS AND DISCUSSION

To understand structure transition of  $\text{BaCO}_3$ , temperature-dependent XRD spectra of  $\text{BaCO}_3$  powder were measured in 1 atm  $\text{CO}_2$ , as shown in Fig. 1(a). The strongest  $2\theta$ -peak appears at  $24.0^\circ$  at room temperature, indicating an orthorhombic phase.<sup>25</sup> Near  $810\text{ }^\circ\text{C}$ , two new peaks appear at  $21.2^\circ$  and  $25.9^\circ$ , suggesting a new phase. As temperature increases, the orthorhombic phase disappears completely near  $850\text{ }^\circ\text{C}$ . These phenomena reveal a first-order orthorhombic-hexagonal<sup>26,27</sup> structural transition in the region of  $810\text{--}850\text{ }^\circ\text{C}$ .

To identify possible decomposition, the XRD spectra of  $\text{ZrO}_2$ ,  $\text{CeO}_2$ , and  $\text{Y}_2\text{O}_3$  powders were obtained at room temperature as given in Fig. 1(b). The main  $2\theta$  peaks appear at  $28.2^\circ$  and  $31.5^\circ$  for monoclinic  $\text{ZrO}_2$  and  $28.5^\circ$ ,  $33.0^\circ$ ,  $47.5^\circ$ , and  $56.2^\circ$  for cubic  $\text{CeO}_2$ . For cubic  $\text{Y}_2\text{O}_3$ , the main  $2\theta$  peaks occur at  $29.2^\circ$ ,  $33.9^\circ$ ,  $48.5^\circ$ , and  $57.8^\circ$ . There is no phase transformation for both  $\text{ZrO}_2$  and  $\text{CeO}_2$  below  $1000\text{ }^\circ\text{C}$ .<sup>28,29</sup>  $\text{BaCeO}_3$  has a phase sequence of orthorhombic ( $Pnma$ )-orthorhombic ( $Imma$ )-rhombohedral ( $R\bar{3}c$ )-cubic ( $Pm\bar{3}m$ ) at 290, 400, and  $900\text{ }^\circ\text{C}$ .<sup>30</sup>  $\text{BaZrO}_3$  is cubic at and above room temperature.<sup>31</sup>

Figure 2(a) shows the XRD spectrum of BZCY442 powder before exposure to  $\text{CO}_2$ . The main  $2\theta$  peaks include (110), (200), (211), (220), (310), and (222) and suggest a

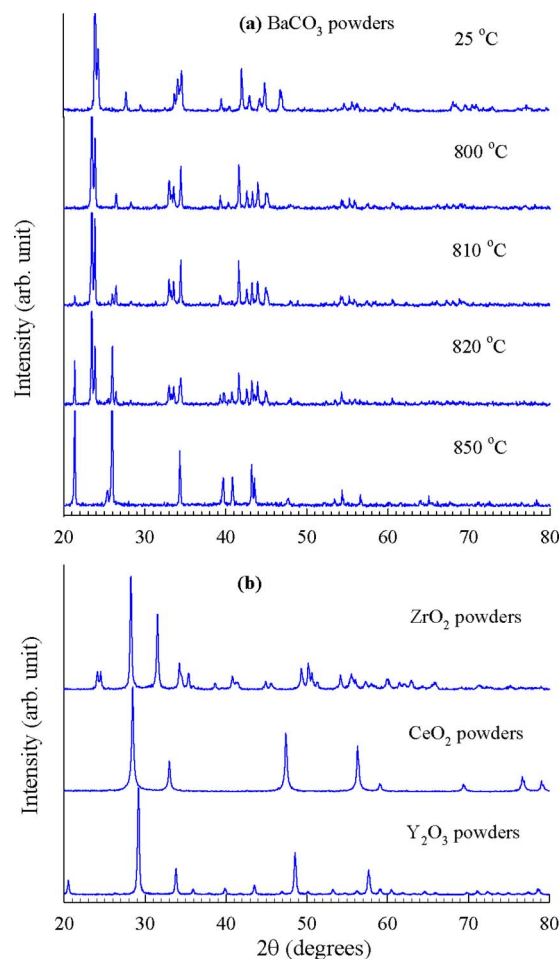


FIG. 1. (Color online) (a) *In situ* XRD spectra of  $\text{BaCO}_3$  powder (99.9%) in 1 atm  $\text{CO}_2$  and (b) XRD spectra of  $\text{ZrO}_2$ ,  $\text{CeO}_2$ , and  $\text{Y}_2\text{O}_3$  powders (99.9%) taken at room temperature.

body-centered cubic structure, but the weak (111) peak is consistent with the expected simple-cubic (sc) perovskite structure, according to a structure-factor calculation.<sup>25</sup> The (100) peak expected for sc is too weak to see in BZCY442 because the reflections from the (100) planes containing BaO are almost exactly cancelled by the reflections from the (100) planes containing  $\text{Zr}_{0.4}\text{Ce}_{0.4}\text{Y}_{0.2}\text{O}_2$ . The cubic lattice parameter was estimated from the strongest (110) peak to be  $a = 4.322\text{ \AA}$ . Two weak peaks as indicated by \* near the (200) and (310) peaks correlate to a second phase.

As shown in Figs. 2(b)–2(e), the XRD spectra remain almost the same below  $550\text{ }^\circ\text{C}$ . At  $600\text{ }^\circ\text{C}$  the orthorhombic peak ( $2\theta=23.7^\circ$ ) of  $\text{BaCO}_3$  appears as denoted by “A.” In addition, four weak broad peaks were observed at  $28.5^\circ$ ,  $33.5^\circ$ ,  $47.9^\circ$ , and  $56.5^\circ$  as indicated by “B.” These weak broad peaks most likely correspond to a  $\text{CeO}_2$ -like structure because their XRD peaks close to main peaks of  $\text{CeO}_2$  [Fig. 1(b)]. The intensities of A and B peaks grow gradually as temperature increases. At  $900\text{ }^\circ\text{C}$  [Fig. 2(j)] A peak disappears and two new peaks appear at  $21.4^\circ$  and  $25.9^\circ$  as indicated by “C,” confirming an orthorhombic-hexagonal transformation of  $\text{BaCO}_3$ .

Figures 3(a)–3(d) show Raman spectra of  $\text{ZrO}_2$ ,  $\text{CeO}_2$ ,  $\text{Y}_2\text{O}_3$ , and  $\text{BaCO}_3$  powders taken at room temperature. The major vibrations are  $474, 461,$  and  $375\text{ cm}^{-1}$  for  $\text{ZrO}_2$ ,

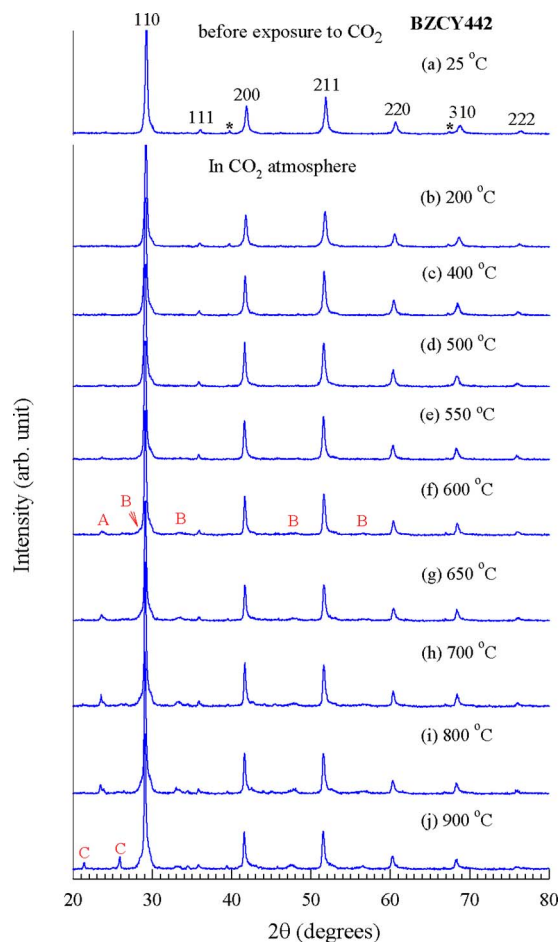


FIG. 2. (Color online) Temperature-dependent XRD spectra of BZCY442. A, B, and C peaks correspond to orthorhombic  $\text{BaCO}_3$ , cubic  $(\text{Zr}, \text{Ce}, \text{Y})\text{O}_2$ , and hexagonal  $\text{BaCO}_3$ , respectively.

$\text{CeO}_2$ , and  $\text{Y}_2\text{O}_3$ , respectively. The  $474\text{ cm}^{-1}$  of  $\text{ZrO}_2$  corresponds to the  $A_g$  Raman-active mode of O–O vibration.<sup>32</sup> The  $461\text{ cm}^{-1}$  vibration of  $\text{CeO}_2$  results from the  $F_{2g}$  Raman-active mode of fluorite structure.<sup>33</sup> The main vibrations of  $\text{BaCO}_3$  include  $690$ ,  $1059$ , and  $1419\text{ cm}^{-1}$  [Fig. 3(d)], which correspond respectively to the doubly degenerate bending mode ( $E''$ ), symmetric stretching vibration ( $A'_1$ ), and doubly degenerate asymmetric stretching vibration ( $E'$ ) of  $\text{CO}_3^{2-}$  ion.<sup>34</sup> The chemical decompositions occurred in BZCY442 were confirmed by  $689$ ,  $1057$ , and  $1418\text{ cm}^{-1}$  for  $\text{BaCO}_3$  and  $466\text{ cm}^{-1}$  for  $(\text{Zr}, \text{Ce}, \text{Y})\text{O}_2$  [Fig. 3(e)]. Note that the  $466\text{ cm}^{-1}$  is between  $461\text{ cm}^{-1}$  of  $\text{CeO}_2$  and  $474\text{ cm}^{-1}$  of  $\text{ZrO}_2$ , but there are no peaks corresponding to the other  $\text{ZrO}_2$  peaks. Another evidence for cubic fluorite structure is that the major Raman vibrations of both  $\text{Zr–Ce–O}$  ( $\text{Ce}/\text{Zr}=1$ ) and Si-doped  $\text{Zr–Ce–O}$  ( $\text{Ce}/\text{Zr}=1$ ) solid solutions appear near  $470\text{ cm}^{-1}$ .<sup>35</sup> It was concluded that Si dopant hardly changes the phase of the solid solution. As shown in Fig. 3(e), the relative intensities of  $1057\text{ cm}^{-1}$  ( $\text{BaCO}_3$ ) and  $466\text{ cm}^{-1}$  [ $(\text{Zr}, \text{Ce}, \text{Y})\text{O}_2$ ] increase significantly after exposure to  $\text{CO}_2$ , but the intensity of  $356\text{ cm}^{-1}$  vibration decreases.

Figure 4(a) shows the XRD spectrum of BZCY532 powder before exposure to  $\text{CO}_2$ . The cubic lattice parameter was estimated to be  $a=4.307\text{ \AA}$ . Some minor shoulders (or split-

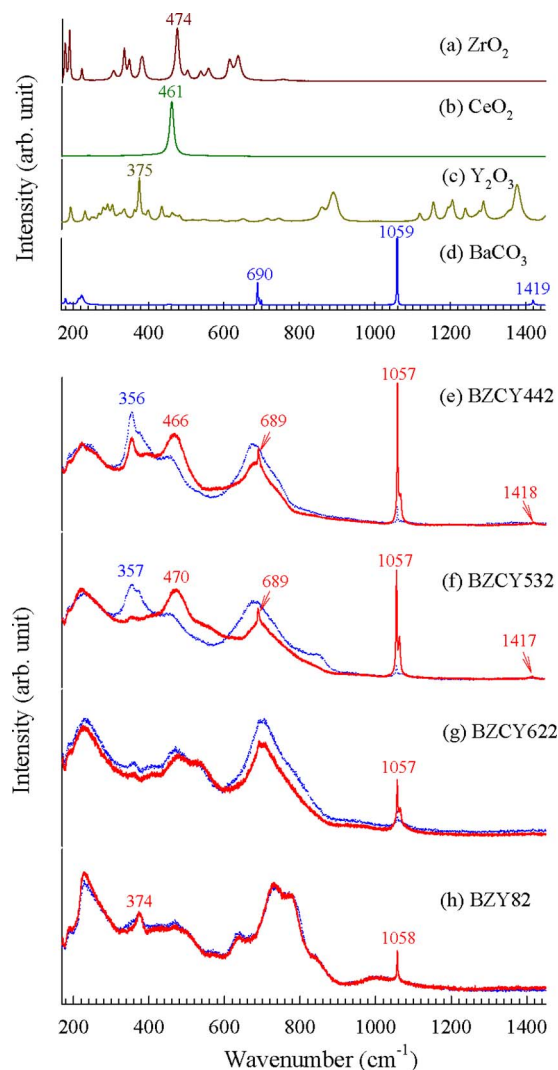


FIG. 3. (Color online) Raman spectra of (a)  $\text{ZrO}_2$ , (b)  $\text{CeO}_2$ , (c)  $\text{Y}_2\text{O}_3$ , (d)  $\text{BaCO}_3$ , (e) BZCY442, (f) BZCY532, (g) BZCY622, and (h) BZY82 powders measured at room temperature. Blue dotted lines are spectra taken before exposure to  $\text{CO}_2$ . Red solid lines are spectra taken after exposure to  $\text{CO}_2$ .

tings) as indicated by \* are possibly caused by a slight rhombohedral distortion from the cubic structure because  $\text{BaZr}_{1-x}\text{Ce}_x\text{O}_3$  is possibly rhombohedral for  $0.15 < x < 0.5$ .<sup>31</sup> In that work, the composition we would call BZCY640 (no Y) was rhombohedral below and cubic above  $477\text{ }^\circ\text{C}$ . Regarding the effect of adding 0.2 fraction Y to replace 0.1 each of Zr and Ce, the closest comparison available is adding 0.1 Y to  $\text{BaCeO}_3$ . This changes the  $Pnma/Imma/R\bar{3}c/c/Pm\bar{3}m$  transition temperatures from  $290/400/900$  to  $237/462/767\text{ }^\circ\text{C}$ .<sup>30</sup> These small but significant changes indicate that a small part of the BZCY532 sample could have been in rhombohedral phase over observed temperature range. This type of perovskite can crystallize in various phases, depending on processing method and sintering temperature. For example, Melekh *et al.*<sup>36</sup> reported that different workers found orthorhombic, cubic, pseudocubic, and tetragonal phases at room temperature for  $\text{BaCeO}_3$  prepared by various methods.

The XRD spectra of BZCY532 powder remain un-

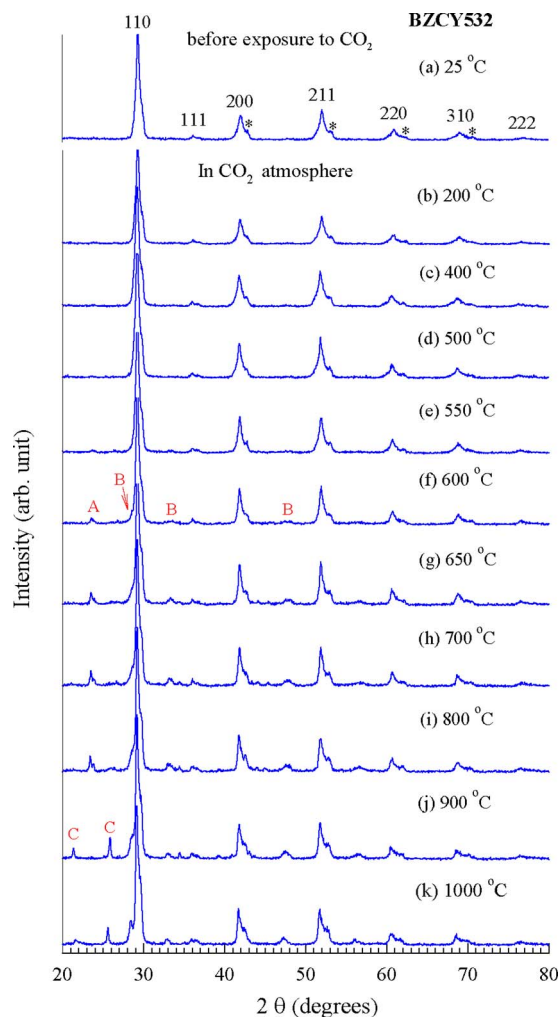


FIG. 4. (Color online) Temperature-dependent XRD spectra of BZCY532.

changed below 550 °C (Fig. 4).  $\text{BaCO}_3$  and  $(\text{Zr,Ce,Y})\text{O}_2$  structures appear noticeably at 550 and 600 °C as indicated by A and B, respectively. They increase in amplitude with temperature and the A peak splits above 800 °C because of the orthorhombic-hexagonal transition of  $\text{BaCO}_3$ . As shown in Fig. 3(f), the decomposition of BZCY532 was confirmed by the 470  $\text{cm}^{-1}$  vibration of  $(\text{Zr,Ce,Y})\text{O}_2$  and 689, 1057, and 1417  $\text{cm}^{-1}$  vibrations of  $\text{BaCO}_3$ . It was found that BZCY532 can absorb about 0.03 g of  $\text{CO}_2$  per ceramic gram after exposure to  $\text{CO}_2$  from 900 °C.

The XRD spectrum of BZCY622 powder before exposure to  $\text{CO}_2$  is given in Fig. 5(a). The cubic lattice parameter was estimated to be  $a=4.279$  Å. Two very weak peaks as indicated by \* near the (200) and (310) peaks correlate to a second phase. As shown in Fig. 5, BZCY622 remains thermally stable in  $\text{CO}_2$  up to 900 °C with a very slight decomposition above 550 K [Figs. 5(f)–5(j)]. This stability is confirmed by the Raman spectra [Fig. 3(g)], in which there is no much difference before and after exposure to  $\text{CO}_2$ . The weak 1057  $\text{cm}^{-1}$  vibration before exposure to  $\text{CO}_2$  indicates a very small amount of  $\text{BaCO}_3$  likely formed in the calcining process.

The Raman spectra of BZCY442, BZCY532, and BZCY622 powders before exposure to  $\text{CO}_2$  [Figs. 3(e)–3(g)]

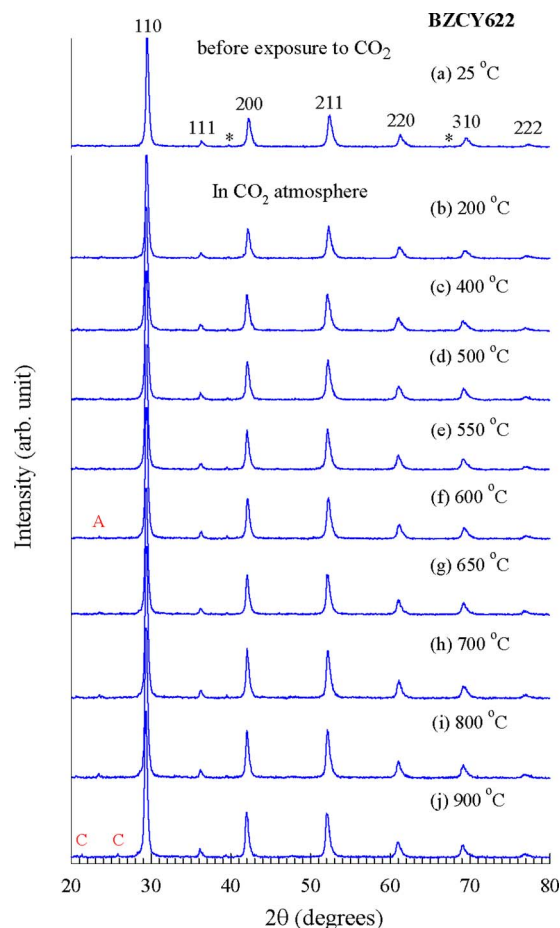


FIG. 5. (Color online) Temperature-dependent XRD spectra of BZCY622.

show that the 356 (or 357  $\text{cm}^{-1}$ ) vibrations and its shoulder (about 370  $\text{cm}^{-1}$ ) become weaker as cerium content decreases, indicating that the 356 (or 357) and 370  $\text{cm}^{-1}$  vibrations are sensitive to cerium content. These are consistent with the Raman result of  $\text{BaZr}_{1-x}\text{Ce}_x\text{O}_3$ ,<sup>31</sup> in which the relative intensities of these vibrations (356 and 370  $\text{cm}^{-1}$ ) decrease and their frequencies also shift to higher region as  $x$  decreases.

Just as for BZCY622, BZY82 powder demonstrates a good thermal stability in  $\text{CO}_2$  as evidenced in Figs. 6 and 3(h). Some weak peaks of  $\text{Y}_2\text{O}_3$ -like structure were observed as indicated by “D,” which was confirmed by the Raman vibration of 374  $\text{cm}^{-1}$  [Figs. 3(c) and 3(h)]. The cubic lattice parameter of BZY82 is  $a=4.216$  Å. This is slightly larger than the lattice constant  $a=4.192$  Å of cubic  $\text{BaZrO}_3$  (Ref. 31) because the ionic radius of  $\text{Y}^{3+}$  ( $R^{\text{III}}=0.9$  Å) is larger than for  $\text{Zr}^{4+}$  ( $R^{\text{IV}}=0.72$  Å). The lattice parameters of BZY82 (4.216 Å), BZCY622 (4.279 Å), BZCY532 (4.307 Å), and BZCY442 (4.322 Å) increase with Ce content because the radius of  $\text{Ce}^{4+}$  ( $R^{\text{IV}}=0.87$  Å) is larger than for  $\text{Zr}^{4+}$  ( $R^{\text{IV}}=0.72$  Å).<sup>37</sup> Figure 7 shows a nearly linear relation between lattice parameter and Zr content.

As shown in the results and discussion described above, no apparent decompositions of  $\text{BaCO}_3$  and  $(\text{Zr,Ce,Y})\text{O}_2$  were detected below 550 °C, and the BZCY622 and BZY82 exhibit significant thermal stability in  $\text{CO}_2$  at intermediate temperatures (600–900 °C). To clarify the kinetic stability

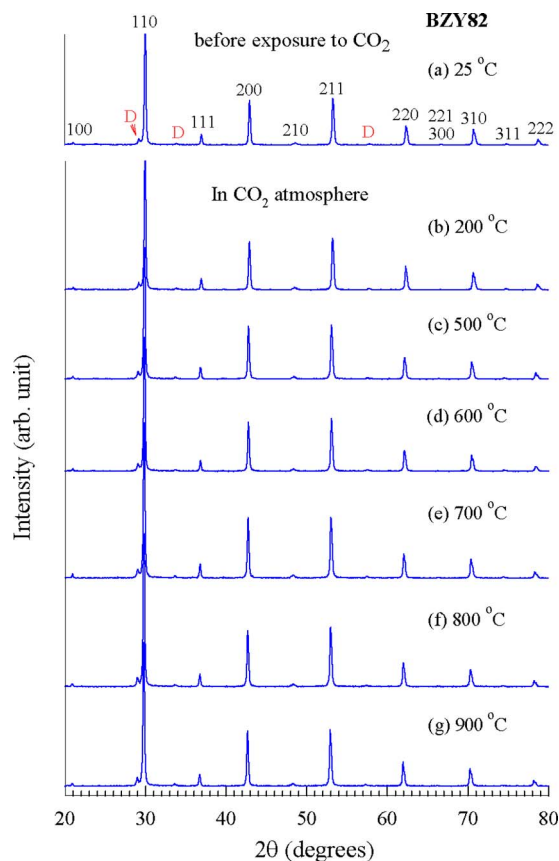


FIG. 6. (Color online) Temperature-dependent XRD spectra of BZY82. D indicates the  $Y_2O_3$ -like structure.

below and above 550 °C, time-dependent XRD spectra in  $CO_2$  were measured at 500 and 700 °C separately, as shown in Fig. 8. The term kinetic stability describes a metastable chemical state that is trapped because of a high activation energy barrier, which causes a very slow reaction rate or essentially none at all, to reach the stable state.

After a long duration (about 29 h) in 1 atm  $CO_2$  at 500 °C [Fig. 8(a)], no apparent decomposition was detected in BZCY442 and BZCY532, indicating a kinetic stability below 500 °C. However, at 700 °C [Fig. 8(b)] both BZCY442 and BZCY532 showed an obvious chemical decomposition as indicated by A and B peaks, which grows as the exposure time increases. Note that both orthorhombic

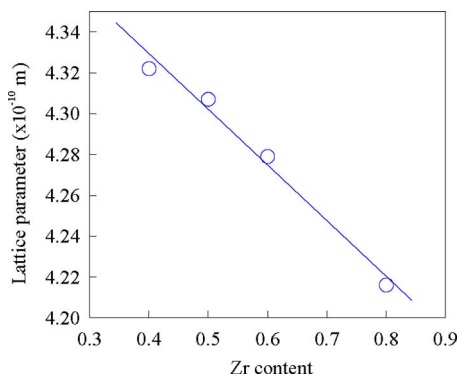


FIG. 7. (Color online) Lattice parameter vs Zr content before exposure to  $CO_2$ .

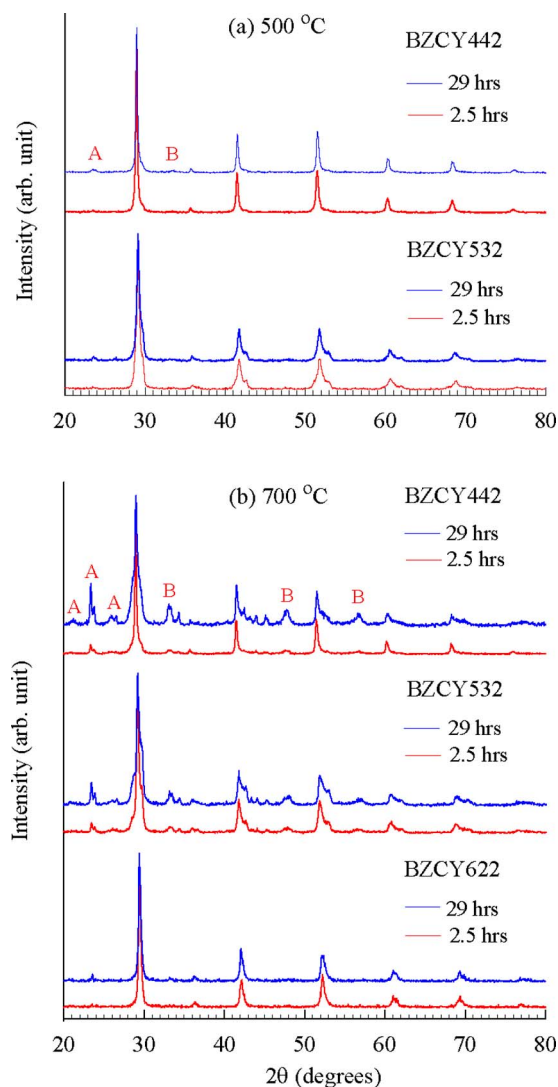


FIG. 8. (Color online) Time-dependent XRD spectra in  $CO_2$  at 500 and 700 °C. A and B correspond to orthorhombic and hexagonal  $BaCO_3$  and cubic  $(Zr,Ce,Y)O_2$ , respectively.

and hexagonal structures of  $BaCO_3$  appear after a long duration at 700 °C, indicating that the orthorhombic-hexagonal structural transformation could occur at lower temperature if the exposure time in  $CO_2$  is longer enough. The time-dependent relative intensity between the A peak ( $2\theta \cong 24^\circ$ ) of  $BaCO_3$  and the (110) peak ( $2\theta \cong 29^\circ$ ) is plotted in Fig. 9. The relative intensity of  $BaCO_3$  at 700 °C grows obviously as exposure time increases in both BZCY442 and BZCY532, especially for BZCY442 which exhibits a rapid decomposition of  $BaCO_3$  and  $(Zr,Ce,Y)O_2$ . Compared with BZCY442 and BZCY532, BZCY622 demonstrates a much better kinetic stability at 700 °C. The kinetic instability of BZCY442 and BZCY532 at 700 °C was also confirmed by the post-Raman spectra, as shown in Figs. 10(a) and 10(b), respectively. As seen in Fig. 10(c), no apparent decomposition was detected in BZCY622 after a long duration in  $CO_2$  at 700 °C.

As shown in Fig. 9, both BZCY442 and BZCY532 exhibit similar slopes of relative intensity versus exposure time at 700 °C below 13 h. In BZCY442,  $BaCO_3$  concentration rises almost linearly with exposure time with an extrapolated

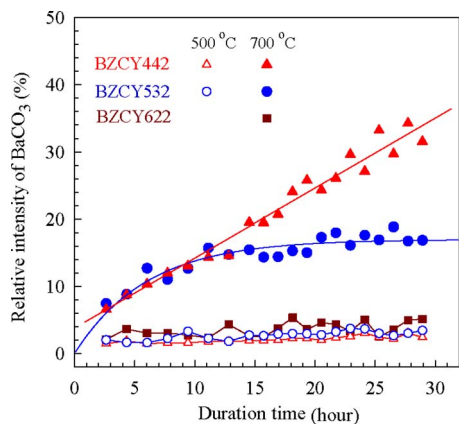


FIG. 9. (Color online) Time-dependent relative intensity between the  $\text{BaCO}_3$  peak ( $2\theta \cong 24^\circ$ ) and (110) peak at 500 and 700 °C. The solid blue curve is the fitting of the exponential rise equation for BZCY532 at 700 °C.

zero-time intercept of 4%. For BZCY532, the time-dependent relative intensity at 700 °C can be approximately described by the exponential rise equation  $n(t) = n_s [1 - e^{-t/\tau}]$  with  $n_s = 17\%$  and  $\tau = 6$  h as indicated by the solid curve.  $n(t)$  is the time-dependent  $\text{BaCO}_3$  concentration and  $\tau$  is the rise time constant. The saturation concentration of  $n_s = 17\%$  implies that the carbonate forms a protective layer that keeps  $\text{CO}_2$  from reacting with the interior of the powders.

#### IV. CONCLUSIONS

$\text{Ba}(\text{Zr}_{0.8-x}\text{Ce}_x\text{Y}_{0.2})\text{O}_{2.9}$  ( $x=0.0, 0.2, 0.3,$  and  $0.4$ ) ceramics exhibit a promising thermal stability below 550 °C with-

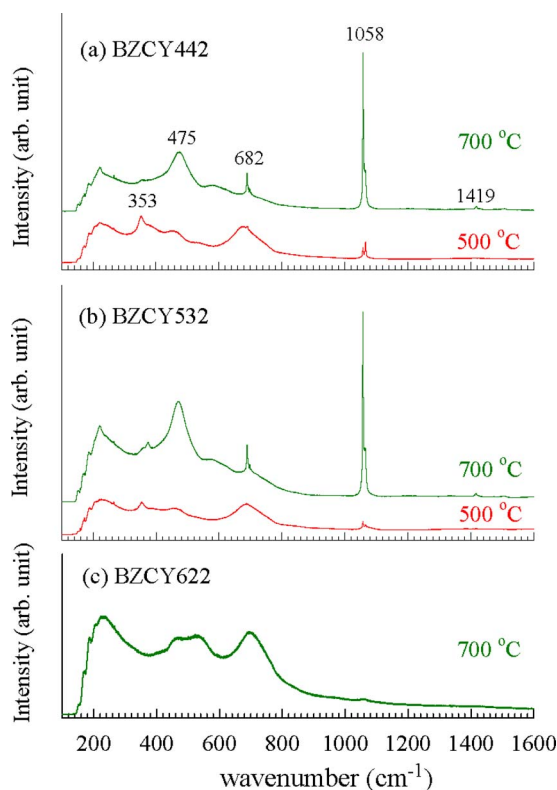


FIG. 10. (Color online) Post-Raman spectra of (a) BZCY442, (b) BZCY532, and (c) BZCY622 powders after long-duration exposures in  $\text{CO}_2$  at 500 and 700 °C.

out apparent chemical decomposition in 1 atm  $\text{CO}_2$ . As temperature rises, the low-cerium compounds  $\text{Ba}(\text{Zr}_{0.8-x}\text{Ce}_x\text{Y}_{0.2})\text{O}_{2.9}$  ( $x \leq 0.2$ ) retain kinetic stability up to 1000 °C after a long-duration exposure to  $\text{CO}_2$ . An orthorhombic-hexagonal structure transformation of  $\text{BaCO}_3$  was evidenced in the region of 810–850 °C in BZCY442 and BZCY532 upon heating. Raman vibrations of 1057 and 466  $\text{cm}^{-1}$  exhibit high sensitivity to appearances of  $\text{BaCO}_3$  and  $(\text{Zr}, \text{Ce}, \text{Y})\text{O}_2$ , respectively. Both XRD spectra and Raman vibrations of  $\text{ZrO}_2$ ,  $\text{CeO}_2$ , and  $\text{Y}_2\text{O}_3$  powders were also obtained at room temperature and can be useful for identification of these compounds. This study suggests that  $\text{Ba}(\text{Zr}_{0.8-x}\text{Ce}_x\text{Y}_{0.2})\text{O}_{2.9}$  ceramics with  $x \leq 0.2$  are promising candidates for proton-conducting SOFC applications at intermediate temperatures (700–850 °C).

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- <sup>1</sup>J. P. Tremblay, R. S. Gemmen, and D. J. Bayless, *J. Power Sources* **163**, 986 (2007).
- <sup>2</sup>R. S. Gemmen and J. P. Tremblay, *J. Power Sources* **161**, 1084 (2006).
- <sup>3</sup>M. Gong, X. Liu, J. P. Tremblay, and C. Johnson, *J. Power Sources* **168**, 289 (2007).
- <sup>4</sup>S. McIntosh and R. J. Gorte, *Chem. Rev. (Washington, D.C.)* **104**, 4845 (2004).
- <sup>5</sup>S. M. Haile, *Acta Mater.* **51**, 5981 (2003).
- <sup>6</sup>C. W. Tanner and A. V. Virkar, *J. Electrochem. Soc.* **143**, 1386 (1996).
- <sup>7</sup>S. V. Bhide and A. V. Virkar, *J. Electrochem. Soc.* **146**, 2038 (1999).
- <sup>8</sup>N. Zakowsky, S. Williamson, and J. T. S. Irvine, *Solid State Ionics* **176**, 3019 (2005).
- <sup>9</sup>B. R. Sneha and V. Thangadurai, *J. Solid State Chem.* **180**, 2661 (2007).
- <sup>10</sup>F. Chen, O. T. Sorensen, G. Meng, and D. Peng, *J. Mater. Chem.* **7**, 481 (1997).
- <sup>11</sup>G. Ma, T. Shimura, and H. Iwahara, *Solid State Ionics* **110**, 103 (1998).
- <sup>12</sup>K. Katahira, Y. Kohchi, T. Shimura, and H. Iwahara, *Solid State Ionics* **138**, 91 (2000).
- <sup>13</sup>Z. Zhong, *Solid State Ionics* **178**, 213 (2007).
- <sup>14</sup>F. M. M. Snijkers, A. Buekenhoudt, J. Coymans, and J. J. Luyten, *Scr. Mater.* **50**, 655 (2004).
- <sup>15</sup>A. Magrez and T. Schober, *Solid State Ionics* **175**, 585 (2004).
- <sup>16</sup>K. H. Ryu and S. M. Haile, *Solid State Ionics* **125**, 355 (1999).
- <sup>17</sup>S. M. Haile, G. Staneff, and K. H. Ryu, *J. Mater. Sci.* **36**, 1149 (2001).
- <sup>18</sup>C. Zuo, S. Zha, M. Liu, M. Hatano, and M. Uchiyama, *Adv. Mater. (Weinheim, Ger.)* **18**, 3318 (2006).
- <sup>19</sup>A. S. Patnaik and A. V. Virkar, *J. Electrochem. Soc.* **153**, A1397 (2006).
- <sup>20</sup>P. Babilo, T. Uda, and S. M. Haile, *J. Mater. Res.* **22**, 1322 (2007).
- <sup>21</sup>P. Babilo and S. M. Haile, *J. Am. Ceram. Soc.* **88**, 2362 (2005).
- <sup>22</sup>A. K. Azad and J. T. S. Irvine, *Solid State Ionics* **178**, 635 (2007).
- <sup>23</sup>S. S. Bhella and V. Thangadurai, *J. Power Sources* **186**, 311 (2009).
- <sup>24</sup>L. A. Chick, L. R. Pederson, G. D. Maupin, J. L. Bates, L. E. Thomas, and G. J. Exarhos, *Mater. Lett.* **10**, 6 (1990).
- <sup>25</sup>B. D. Cullity, *Elements of X-ray Diffraction* (Addison-Wesley, Reading, MA, 1978).
- <sup>26</sup>R. Strobel, M. Maciejewski, S. E. Pratsinis, and A. Baiker, *Thermochim. Acta* **445**, 23 (2006).
- <sup>27</sup>C. R. M. Rao and P. N. Mehrotra, *J. Therm. Anal.* **17**, 539 (1979).
- <sup>28</sup>M. Yoshimura, *Am. Ceram. Soc. Bull.* **67**, 1950 (1988).
- <sup>29</sup>M. Yashima, D. Ishimura, Y. Yamaguchi, K. Ohoyama, and K. Kawachi, *Chem. Phys. Lett.* **372**, 784 (2003).
- <sup>30</sup>K. S. Knight, *Solid State Ionics* **145**, 275 (2001).
- <sup>31</sup>I. Charrier-Cougoulic, T. Pagnier, and G. Lucazeau, *J. Solid State Chem.* **142**, 220 (1999).
- <sup>32</sup>B.-K. Kim and H.-O. Hamaguchi, *Phys. Status Solidi B* **203**, 557 (1997).
- <sup>33</sup>R. Q. Long, Y. P. Huang, and H. L. Wan, *J. Raman Spectrosc.* **28**, 29

(1997).

<sup>34</sup>R. L. Frost and J. M. Bouzaid, *J. Raman Spectrosc.* **38**, 873 (2007).

<sup>35</sup>X. Wang, G. Lu, Y. Guo, Y. Xue, L. Jiang, Y. Guo, and Z. Zhang, *Catal. Today* **126**, 412 (2007).

<sup>36</sup>B.-T. Melekh, V. M. Egorov, Y. M. Baikov, N. F. Kartenko, Y. N. Filin, M. E. Kompan, I. I. Novak, G. B. Venus, and V. B. Kulik, *Solid State Ionics* **97**, 465 (1997).

<sup>37</sup>R. D. Shannon, *Acta Crystallogr.* **A32**, 751 (1976).